

Fixed points of transformation of a renormalization group in a generalized fermionic hierarchical model

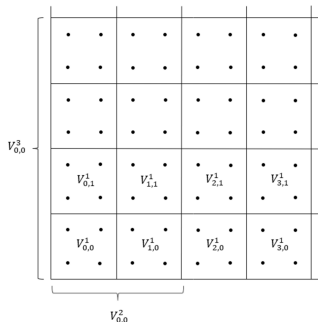
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Slide 1 Renormalization group in a generalized hierarchical model

$$\begin{array}{ccc}
 (Z^d, R) & \longrightarrow & (Z^d, \Gamma) \\
 \downarrow & & \downarrow \\
 (T^d, R) & \longrightarrow & (T^d, \Gamma)
 \end{array}$$



$$\phi'(i) \equiv (r(\alpha)\phi)(i) = 2^{-\alpha/2} \sum_{j \in V_i^1} \phi(j)$$

Slide 2 Renormalization group in a generalized hierarchical model

Let $T = \{0, 1, 2, \dots\}$ and $V_k^s = \{j \in T : k \cdot 2^s \leq j < (k+1) \cdot 2^s\}$, where $k \in T$, $s \in N = \{1, 2, 3, \dots\}$.

The hierarchical distance $d_2(i, j)$, $i, j \in T$, $i \neq j$ is defined as $d_2(i, j) = 2^{s(i, j)}$, where

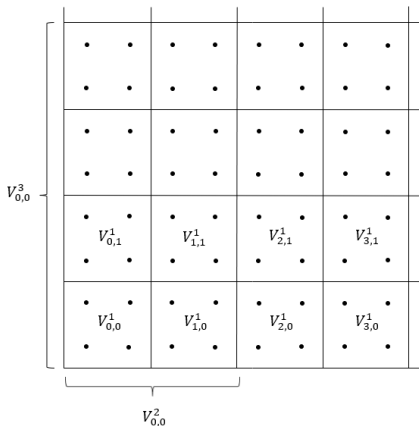
$$s(i, j) = \min\{s : \text{is } k \in T \text{ such that } i, j \in V_k^s\}$$

Let $T^2 = T \times T$, $k = (k_1, k_2) \in T^2$,

$$V_k^s = \{(j_1, j_2) \in T^2 : k_1 \cdot 2^s \leq j_1 < (k_1+1) \cdot 2^s, k_2 \cdot 2^s \leq j_2 < (k_2+1) \cdot 2^s\}.$$

For any $k = (k_1, k_2) \in T^2$, $l = (l_1, l_2) \in T^2$, $k \neq l$ we define $s(k, l) = \max(s(k_1, l_1), s(k_2, l_2))$. The hierarchical distance on T^2 is defined as $d_2(k, l) = 2^{s(k, l)}$.

Slide 3 Renormalization group in a generalized hierarchical model



Slide 4 Renormalization group in a generalized hierarchical model

Let us consider a 4-component fermionic field

$$\psi^*(i) = (\bar{\psi}_1(i), \psi_1(i), \bar{\psi}_2(i), \psi_2(i)), \quad i \in T^2,$$

all components of which are generators of the Grassmann algebra. Let us redenote V_0^N by Λ_N . Let Γ_N be the Grassmann subalgebra generated by all generators

$$\bar{\psi}_1(i), \psi_1(i), \bar{\psi}_2(i), \psi_2(i), i \in \Lambda_N.$$

The block-spin renormalization group transformation is defined by the formula

$$\psi^{*'}(i) \equiv (r(\alpha)\psi^*)(i) = 2^{-\alpha/2} \sum_{j \in V_i^1} \psi^*(j), \quad (1)$$

where $\alpha \in R^1$ is the renormalization group parameter.

Slide 5 Renormalization group in a generalized hierarchical model

The Gaussian fermionic field with zero mean and binary correlation function

$$\langle \psi_n(k) \bar{\psi}_m(l) \rangle = \delta_{n,m} b(k, l), \quad n, m = 1, 2, \quad k, l \in T^2$$

is defined on the entire lattice as a quasi-state (expectation) $\langle \cdot \rangle$ on the algebra of all monomials such that $\langle F \rangle$ for a monomial of even degree F is calculated according to Wick's rules and $\langle F \rangle = 0$ for any monomial of odd degree, $\delta_{n,m}$ is the Kronecker symbol. Here we use the following notations:

$$\psi(i) = (\psi_1(i), \psi_2(i)), \quad \bar{\psi}(i) = (\bar{\psi}_1(i), \bar{\psi}_2(i)),$$

$$\bar{\psi}(i) \eta(i) = \bar{\psi}_1(i) \eta_1(i) + \bar{\psi}_2(i) \eta_2(i), \quad i \in T^2.$$

Slide 6 Renormalization group in a generalized hierarchical model

We define the following functions on T^2 :

$$d(k, l; \lambda) = d_2(k, l), \text{ if } s(k_1, l_1) \neq s(k_2, l_2),$$

$$d(k, l; \lambda) = \lambda d_2(k, l), \text{ if } s(k_1, l_1) = s(k_2, l_2),$$

$$f(k, l; \lambda; \alpha) = d^\alpha(k, l; \lambda), \text{ if } k \neq l, \quad f(k, k; \lambda; \alpha) = \frac{2 + \lambda^\alpha}{4(1 - 2^{-(2+\alpha)})},$$

λ is a real parameter, $\lambda > 0$.

Slide 7 Renormalization group in a generalized hierarchical model

Let $b(k, l; \lambda; \alpha) = f(k, l; \lambda; \alpha - 4)$. It was shown in [1] that a zero-mean Gaussian fermionic field with a binary correlation function

$$\langle \psi_n(k) \bar{\psi}_m(l) \rangle = \delta_{n,m} b(k, l; \lambda; \alpha), \quad n, m = 1, 2, \quad k, l \in T^2$$

is invariant under the transformation of the renormalization group with the parameter α (1). Let us denote the corresponding Gaussian quasi-state as $\rho_0(\lambda; \alpha)$.

Slide 8 Renormalization group in a generalized hierarchical model

Consider the restriction of the Gaussian state $\rho_0(\lambda; \alpha)$ to the volume Γ_N . It follows from [1, Theorem 1] that for any $F(\psi^*) \in \Gamma_N$:

$$\rho_0(\lambda; \alpha)(F(\psi^*)) = Z_N^{-1}(\lambda; \alpha) \int F(\psi^*) \exp\{-H_{0,N}(\psi^*; \lambda; \alpha)\} d\psi^*,$$

$$d\psi^* = \prod_{i \in \Lambda_N} d\psi_1(i) d\bar{\psi}_1(i) d\psi_2(i) d\bar{\psi}_2(i),$$

integration is carried out according to the rules of superanalysis,

$$Z_N(\lambda; \alpha) = \int \exp\{-H_{0,N}(\psi^*; \lambda; \alpha)\} d\psi^*.$$

Slide 9 Renormalization group in a generalized hierarchical model

Consider the local potential (self-action) for a 4-component fermionic field:

$$L(\psi^*; r, g) = r(\bar{\psi}_1\psi_1 + \bar{\psi}_2\psi_2) + g\bar{\psi}_1\psi_1\bar{\psi}_2\psi_2.$$

Denote $u(\psi^*(i)) = \exp\{-L(\psi^*(i); r, g)\}$. We will use the notation $\rho_N(\lambda; \alpha; u)$ for the state in the algebra Γ_N defined by the formula:

$$\rho_N(\lambda; \alpha; u)(F) = Z_N^{-1}(\lambda; \alpha; u) \int F \exp\{-H_{0,N}(\psi^*; \lambda; \alpha)\} \prod_{i \in \Lambda_N} u(\psi^*(i)) d\psi^*,$$

$$Z_N(\lambda; \alpha; u) = \int \exp\{-H_{0,N}(\psi^*; \lambda; \alpha)\} \prod_{i \in \Lambda_N} u(\psi^*(i)) d\psi^*.$$

Slide 10 Renormalization group in a generalized hierarchical model

If ρ is a state on Γ_N , then the renormalized state of ρ' is defined on Γ_{N-1} according to the formula: $\rho' (F(\psi^*)) = \rho (F(r(\alpha)\psi^*))$. Let's define a quadratic form:

$$\sum_{i,j \in V_k^1} f(i,j;\mu(\lambda);-\alpha) \bar{\eta}^*(i) \eta^*(j) = Q_k(\eta^*), \quad k \in \Lambda_{N-1}.$$

$$f(i,j;\mu;-\alpha) = \frac{2 + \mu^{-\alpha}}{4(1 - 2^{\alpha-2})}, \quad \text{if } i = j,$$

$$\mu(\lambda) = \left(\frac{4(1 - 2^{2-\alpha}) - 3\lambda^{\alpha-4}}{\lambda^{\alpha-4}(1 - 2^{2-\alpha}) - 3 \cdot 2^{2-\alpha}} \right)^{-\frac{1}{\alpha}},$$

$$f(i,j;\mu;-\alpha) = 2^{-\alpha}, \quad \text{if } d_2(i,j) = 1 \text{ and } s(i_1,j_1) \neq s(i_2,j_2), \quad (2)$$

$$f(i,j;\mu;-\alpha) = (2\mu)^{-\alpha}, \quad \text{if } d_2(i,j) = 1 \text{ and } s(i_1,j_1) = s(i_2,j_2).$$

Slide 11 Renormalization group in a generalized hierarchical model

In [1, Theorem 2], it was proved that

$$\rho'_N(\lambda; \alpha; u) = \rho_{N-1}(\lambda; \alpha; u')$$

where

$$u'(\psi^{*'}) = \int \delta \left(\sum_{i \in V_0^1} \eta^*(i) \right) \exp\{-Q_0(\eta^*)\} \quad * \\ * \prod_{i \in V_0^1} \left(u(2^{\alpha/2-2}\psi^{*'} + \eta^*(i)) \right) d\eta^*(i). \quad (3)$$

Let us denote the transformation in the density space given by the right side (3) as $R(\alpha)u$. From this formula we see that the transformation of the renormalization group $R(\alpha)$ is local and does not depend on N .

Slide 12 Renormalization group in a generalized hierarchical model

In this section, instead of regular densities of the form

$$u(\psi^*) = \exp\{-L(\psi^*; r, g)\}$$

we will use Grassmann-valued densities of the general form

$$u(\psi^*) = c_0 + c_1(\bar{\psi}_1\psi_1 + \bar{\psi}_2\psi_2) + c_2\bar{\psi}_1\psi_1\bar{\psi}_2\psi_2.$$

Slide 13 Renormalization group in a generalized hierarchical model

The triple $c = (c_0, c_1, c_2)$ can be naturally treated as a point in the two-dimensional real projective space because two sets that differ by a nonzero factor represent the same Gibbs state. Let

$$\delta = 2 \frac{(2^{2-\alpha} - 1)(\lambda^{\alpha-4} - 1)}{\lambda^{\alpha-4} + 2^{4-\alpha} - 2} + 2, \gamma = 2^{\alpha-2}.$$

Slide 14 Renormalization group in a generalized hierarchical model

Theorem

Let $u'(\psi^{*'}) = R(\alpha)u$. Then

$$u'(\psi^{*'}) = c'_0 + c'_1(\bar{\psi}'_1\psi'_1 + \bar{\psi}'_2\psi'_2) + c'_2\bar{\psi}'_1\psi'_1\bar{\psi}'_2\psi'_2 \quad (4)$$

$$\begin{aligned} c'_0 = f_0(c; \delta) = & c_0^4(\delta^2 - 2\delta + 1) + c_0^3c_1(-4\delta^2 + 6\delta - 2) + \\ & + c_0^3c_2(\delta^2 - \delta + \frac{1}{4}) + c_0^2c_1^2(5\delta^2 - 3\delta - \frac{5}{4}) + c_0^2c_1c_2(-2\delta^2 - \delta + 1) + \\ & + c_0^2c_2^2(\frac{\delta^2}{4} + \frac{1}{2}) + c_0c_1^3(-2\delta^2 - 3\delta + 3) + c_0c_1^2c_2(\frac{\delta^2}{2} + 5\delta - \frac{3}{2}) + \\ & + c_0c_1c_2^2(-\delta - 1) + \frac{c_0c_2^3}{4} + c_1^4(\frac{\delta^2}{4} + \delta) + c_1^3c_2(-\delta - 1) + \frac{3c_1^2c_2^2}{4}, \quad (5) \end{aligned}$$

Slide 15 Renormalization group in a generalized hierarchical model

Theorem

$$\begin{aligned} c_1' = \gamma f_1(c; \delta) = & \gamma(c_0^3 c_1(\delta^2 - 2\delta + 1) + c_0^3 c_2(-\frac{\delta^2}{2} + \frac{3\delta}{4} - \frac{1}{4}) + \\ & + c_0^2 c_1^2(-\frac{7\delta^2}{2} + \frac{21\delta}{4} - \frac{7}{4}) + c_0^2 c_1 c_2(\frac{5\delta^2}{2} - 2\delta) + c_0^2 c_2^2(-\frac{\delta^2}{4} - \frac{\delta}{2} + \frac{1}{2}) + \\ & + c_0 c_1^3(\frac{7\delta^2}{2} - 2\delta - 1) + c_0 c_1^2 c_2(-3\delta^2 - \frac{3\delta}{2} + \frac{3}{2}) + c_0 c_1 c_2^2(\frac{\delta^2}{2} + 2\delta) + \\ & + c_0 c_2^3(-\frac{\delta}{4} - \frac{1}{4}) + c_1^4(-\frac{3\delta^2}{4} - 2\delta + 2) + c_1^3 c_2(\frac{\delta^2}{2} + 4\delta - 1) + \\ & + c_1^2 c_2^2(-\frac{7\delta}{4} - \frac{7}{4}) + c_1 c_2^3) \quad (6) \end{aligned}$$

Slide 16 Renormalization group in a generalized hierarchical model

Theorem

$$\begin{aligned} c_2' = \gamma^2 f_2(c; \delta) = & \gamma^2 (c_0^3 c_2 (\frac{\delta^2}{4} - \frac{\delta}{2} + \frac{1}{4}) + c_0^2 c_1^2 (\frac{3\delta^2}{4} - \frac{3\delta}{2} + \frac{3}{4}) + \\ & + c_0^2 c_1 c_2 (-2\delta^2 + 3\delta - 1) + c_0^2 c_2^2 (\frac{3\delta^2}{4} - \delta + \frac{1}{2}) + c_0 c_1^3 (-2\delta^2 + 3\delta - 1) + \\ & + c_0 c_1^2 c_2 (4\delta^2 - 2\delta - \frac{3}{2}) + c_0 c_1 c_2^2 (-2\delta^2 - \delta + 1) + c_0 c_2^3 (\frac{\delta^2}{4} + \frac{\delta}{2} + \frac{1}{4}) + \\ & + c_1^4 (\frac{5\delta^2}{4} - \delta) + c_1^3 c_2 (-2\delta^2 - 3\delta + 3) + c_1^2 c_2^2 (\frac{3\delta^2}{4} + \frac{11\delta}{2} - \frac{5}{4}) + \\ & + c_1 c_2^3 (-2\delta - 2) + c_2^4) \quad (7) \end{aligned}$$

Slide 1 Fixed points

Note that the mapping given by the formulas (5), (6), (7) depends not only on α , but also on the parameter δ . Therefore, let us denote the mapping of the renormalization group $R(\alpha)$ in c -space as $R(\alpha; \delta)$:

$$R(\alpha; \delta)(c_0, c_1, c_2) = (f_0(c; \delta), f_1(c; \delta), f_2(c; \delta)).$$

Slide 2 Fixed points

Next, consider the representation of our model in (r, g) -coordinates.
Let:

$$u(\psi^*) = \exp\{-r(\bar{\psi}_1\psi_1 + \bar{\psi}_2\psi_2) + g\bar{\psi}_1\psi_1\bar{\psi}_2\psi_2\}$$

In the projective representation, such a density is given by the triple of coefficients

$$c = (1, -r, r^2 - g).$$

Slide 3 Fixed points

The transformation of the renormalization group in the plane of coupling constants (r, g) will also be denoted by $R(\alpha; \delta)$:

$$(r', g') = R(\alpha; \delta)(r, g).$$

Lemma

The transformation of the renormalization group $R(\alpha; \delta)$ in the plane (r, g) has the form:

$$r' = 2^{\alpha-2} \frac{r_1(r, g)}{r_2(r, g)}, \quad (8)$$

$$g' = \frac{2^{2(\alpha-2)}}{4} g \left(\frac{r_3(r, g)}{r_2(r, g)} \right)^2, \quad (9)$$

Slide 4 Fixed points

Lemma

$$\begin{aligned} r_1(r, g) = & -g^3 r - g^3 \left(\frac{\delta}{4} + \frac{1}{4} \right) + 3g^2 r^3 - g^2 r^2 \left(-\frac{5\delta}{2} - \frac{5}{2} \right) - g^2 r \left(-\frac{\delta^2}{2} - 2\delta \right) - \\ & -g^2 \left(-\frac{\delta^2}{4} - \frac{\delta}{2} + \frac{1}{2} \right) - 3gr^5 - gr^4 \left(\frac{17\delta}{4} + \frac{17}{4} \right) - gr^3 \left(\frac{3\delta^2}{2} + 8\delta - 1 \right) - \\ & -gr^2 \left(\frac{7\delta^2}{2} + \frac{5\delta}{2} - \frac{5}{2} \right) - gr \left(\frac{5\delta^2}{2} - 2\delta \right) - g \left(\frac{\delta^2}{2} - \frac{3\delta}{4} + \frac{1}{4} \right) + r^7 - r^6 (-2\delta - 2) - \\ & -r^5 (-\delta^2 - 6\delta + 1) - r^4 (-4\delta^2 - 4\delta + 4) - r^3 (-6\delta^2 + 4\delta + 1) - \\ & -r^2 (-4\delta^2 + 6\delta - 2) + r(\delta - 1)^2 \end{aligned}$$

Slide 5 Fixed points

Lemma

$$\begin{aligned}r_2(r, g) = & -\frac{g^3}{4} + \frac{3g^2r^2}{2} + g^2r(\delta+1) + g^2\left(\frac{\delta^2}{4} + \frac{1}{2}\right) - \frac{9gr^4}{4} + gr^3(-3\delta-3) + \\& + gr^2(-\delta^2 - 5\delta + \frac{1}{2}) + gr(-2\delta^2 - \delta + 1) + g(-\delta^2 + \delta - \frac{1}{4}) + r^6 + r^5(2\delta+2) + \\& + r^4(\delta^2 + 6\delta - 1) + r^3(4\delta^2 + 4\delta - 4) + r^2(6\delta^2 - 4\delta - 1) + \\& + r(4\delta^2 - 6\delta + 2) + (\delta - 1)^2 \\r_3(r, g) = & (-g + (r+1)^2)\left(\frac{-\delta^2g}{2} - 2\delta gr + 2\delta r^3 + g^2 - 2gr^2 + r^4 + \right. \\& \left. + r^2(\delta^2 + 2\delta - 2) + r(2\delta^2 - 2\delta) + (\delta - 1)^2\right).\end{aligned}$$

Slide 6 Fixed points

Sorollary

Consequence. The domains $\{(r, g) : g > 0\}$ (upper half-plane), $\{(r, g) : g < 0\}$ (lower half-plane) and $\{(r, g) : g = 0\}$ are invariant domains of the mapping $R(\alpha; \delta)$. The action of $R(\alpha; \delta)$ on the line $\{(r, g) : g = 0\}$ is linear:

$$R(\alpha; \delta)(r, 0) = (2^{\alpha-2}r, 0).$$

Slide 7 Fixed points

Let

$$A_0 = (1, 0, 0), \quad A_1 = (0, 0, 1).$$

The points A_0 and A_1 are fixed points of $R(\alpha, \delta)$ for all α and δ . In coordinates (r, g) , the point A_0 is given as $(0, 0)$, so we call A_0 the Gaussian fixed point. The point A_1 does not belong to the regular plane $\{(r, g)\}$, and it corresponds to the density $u(\psi^*) = \bar{\psi}_1 \psi \bar{\psi}_2 \psi_2$ (Grassmann delta function).

Slide 8 Fixed points

Theorem

- 1 *The fixed point A_0 is repulsive for $\alpha > 3$, saddle point - for $2 < \alpha < 3$, and attractive for $\alpha < 2$.*
- 2 *The fixed point A_1 is attractive for $\alpha > 2$, the saddle point is in the interval $1 < \alpha < 2$, and repulsive for $\alpha < 1$.*

Slide 9 Fixed points

- 1 if $\alpha > 3$, then the fixed point A_0 is repulsive,
- 2 if $2 < \alpha < 3$, then A_0 is a saddle point,
- 3 if $\alpha < 2$, then A_0 is an attracting point,
- 4 the value $\alpha = 3$ is a bifurcation value and therefore a new (non-Gaussian) branch of fixed points must exist in a neighborhood of the point A_0 .

Slide 10 Fixed points

- 1 if $\alpha > 2$, then A_1 is an attracting point,
- 2 if $1 < \alpha < 2$, then the point A_1 is a saddle fixed point,
- 3 if $\alpha < 1$, then the point A_1 is a repulsive point,
- 4 the value $\alpha = 1$ is bifurcational, therefore, in the vicinity of the point A_1 , there must be a new branch of fixed points.

Slide 11 Fixed points

We consider the realization of the projective c space in the form of the hemisphere

$$S = \{(c_0, c_1, c_2) : c_0^2 + c_1^2 + c_2^2 = 1, c_0 \geq 0\},$$

where the opposite points of the boundary circle $c_1^2 + c_2^2 = 1$ are identified. To obtain the flat (two-dimensional) picture, we use the orthogonal projection S on the disk $D = \{(c_1, c_2) : c_1^2 + c_2^2 \leq 1\}$.

Slide 12 Fixed points

The regular point (r, g) then corresponds to $(c_1(r, g), c_2(r, g))$.

$$c_1(r, g) = -\frac{r}{\sqrt{1 + r^2 + (r^2 - g)^2}}, \quad c_2(r, g) = \frac{r^2 - g}{\sqrt{1 + r^2 + (r^2 - g)^2}}.$$

We note that the points $(c_1(r, g), c_2(r, g))$ belong to the interior of the disk D . The trivial fixed point $r = 0, g = 0$ is also represented by the point $(0, 0)$ in the (c_1, c_2) coordinates. The fixed point δ in (c_1, c_2) coordinates is determined by the point $(0, 1)$. The line $g = 0$ in the (c_1, c_2) space is described by the curve $l_0 = \{(c_1(r, 0), c_2(r, 0)); r \in R\}$.

Slide 13 Fixed points

Let $v = r + 1, \tau = \delta - 2$, then

Theorem

$$\begin{aligned} r_2(r, g; \delta) = s_2(v, g; \tau) &= -\frac{g^3}{4} + g^2 \left(\frac{\tau^2}{4} + \tau v + \frac{3v^2}{2} \right) - \\ &\quad - \frac{gv^2(2\tau + 3v)^2}{4} + v^4(\tau + v)^2; \\ r_3(r, g; \delta) = s_3(v, g; \tau) &= (-g + v^2) \left(g^2 - \frac{g(\tau + 2v)^2}{2} + v^2(\tau + v)^2 \right); \\ r_1(r, g; \delta) = s_1(v, g; \tau) &= \frac{1}{4}g\tau(g - v^2)^2 - s_2(v, g; \tau) + s_3(v, g; \tau)v. \end{aligned} \tag{10}$$



Slide 14 Fixed points

Note that the parameters τ, γ depend on α , but we will not explicitly designate this dependence.

Theorem

Let $\alpha \neq 2$, the renormalization group transformation on a generalized hierarchical model in the space of coupling constants (v, g) have 4 branches of fixed points, where

$$v_{\pm}^1 = \gamma - 1 + \tau(\gamma - 2\beta_{\pm}^1), g_{\pm}^1 = \frac{\beta_{\pm}^1 - 1}{\beta_{\pm}^1 - \frac{1}{2}} v_{\pm}^1{}^2,$$

$$\tau = \delta(\alpha, \lambda) - 2 = 2 \frac{(2^{2-\alpha} - 1)(\lambda^{\alpha-4} - 1)}{\lambda^{\alpha-4} + 2^{4-\alpha} - 2},$$

$$\beta_{\pm}^1 = \frac{3(\gamma(\tau + 1)^2 - 1) \pm \sqrt{(\gamma(\tau + 1)^2 - 1)^2 + 8(\tau + 1)^2(\gamma - 1)^2}}{4\tau(\tau + 2)},$$

Slide 15 Fixed points

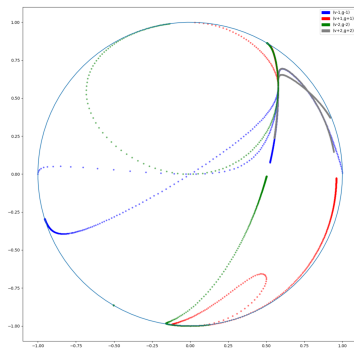
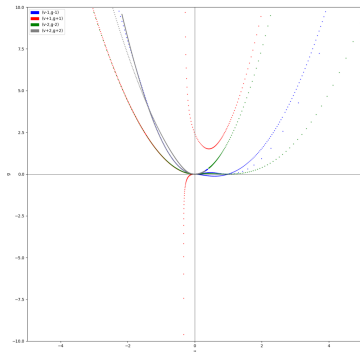
Theorem

$$v_{\pm}^2 = \frac{1-\gamma}{3} - \frac{\tau}{3}(\gamma + 2\beta_{\pm}^2), g_{\pm}^2 = \frac{\beta_{\pm}^2 - 1}{\beta_{\pm}^2 - \frac{1}{2}} v_{\pm}^2,$$

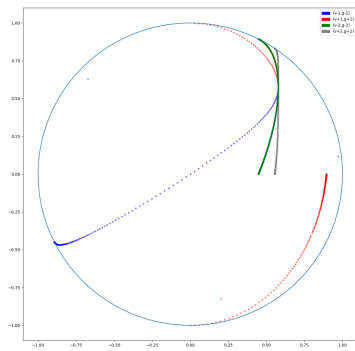
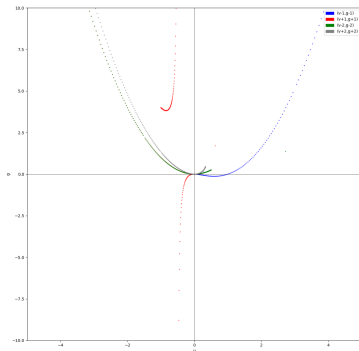
$$\beta_{\pm}^2 = -\frac{5\gamma^2\tau^2 + 2\gamma^2\tau - 11\gamma\tau^2 - 10\gamma\tau + 8\tau - 3(\gamma - 1)^2 \pm 3\sqrt{\Delta(\tau, \gamma)}}{4\tau(-7\gamma\tau - 6\gamma + \tau + 6)},$$

$$\Delta(\tau, \gamma) = \gamma^4(9\tau^4 + 20\tau^3 + 14\tau^2 + 4\tau + 1) + \gamma^3(18\tau^4 + 16\tau^3 - 10\tau^2 - 12\tau - 4) + \gamma^2(9\tau^4 - 20\tau^3 - 14\tau^2 + 12\tau + 6) + \gamma(-16\tau^3 + 2\tau^2 - 4\tau - 4) + 8\tau^2 + 1.$$

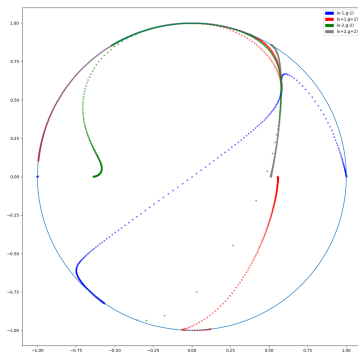
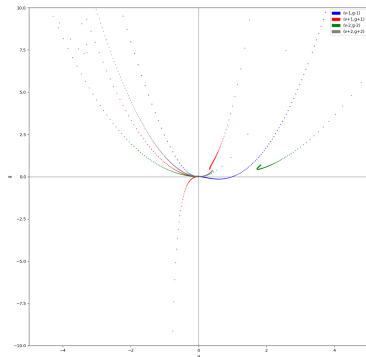
Slide 16 Fig.1 $\lambda = 1/\sqrt{2}$ Fixed points



Slide 17 Fig.2 $\lambda = 1$ Fixed points



Slide 18 Fig.3 $\lambda = \sqrt{2}$ Fixed points



References

- [1] M. D. Missarov , "New Variant of the Fermionic Model on the Hierarchical Two-Dimensional Lattice P-Adic Num Ultrametr Anal Appl, 2020, Vol. 12, pp. 163-170.
- [2] M. D. Missarov, D. A. Khajrullin, "The renormalization group transformation in the generalized fermionic hierarchical model Izv. Math., 2023, Vol. 87, No. 5, pp. 1011-1012.